



# On structural hysteresis damping of vibro-impact dynamics: case study

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- Dynamics of inertial pair coupled only via frictional interface
- Various applications like, for instance, friction drives, items on conveyor belts and turntables, not-flanged manipulation in robotics, etc.



 Well defined solutions, also with discontinuities (i.e. in Filippov sense)



 Less studied dynamics, also with changed cause-effect relationship





• Driving inertial body  $m_1$  and driven inertial body  $m_2$  (in generalized coordinates)



[Ruderman, Zagvozdkin, Rachinskii, IEEE CDC, 2022]

 Coupling between two bodies via the Coulomb friction force

$$f = b \operatorname{sgn}(v), \quad v \equiv \dot{x}_1 - \dot{x}_2$$

b > 0: Coulomb friction coefficient

Three-points signum function

$$\operatorname{sgn}(v) = \begin{cases} 1, & v > 0, \\ 0, & v = 0, \\ -1, & v < 0. \end{cases}$$

⇒ well-defined map for (physical) zero relative velocity



Proposed (new) dynamic model

 $x_1 - x_2 =: z,$  (2)

$$\left(m_1 + m_2 \left(1 - \left|\operatorname{sgn}(\dot{z})\right|\right)\right) \ddot{x}_1 + a_1 \dot{x}_1 + a_2 x_1 + b \operatorname{sgn}(\dot{z}) = u(t), \quad (3)$$

$$m_2 \ddot{x}_1 \left( 1 - \left| \operatorname{sgn}(\dot{z}) \right| \right) \frac{1}{2} \left( 1 - \operatorname{sgn}\left( \left| \ddot{x}_1 \right| - b m_2^{-1} \right) \right) - m_2 \ddot{x}_2 + b \operatorname{sgn}(\dot{z}) = 0.$$
 (4)

• (i) mode:  $m_2$ -body rests on  $m_1$ -body  $(\dot{x}_1 = \dot{x}_2)$ , and no switching  $(|\ddot{x}_1| < bm_2^{-1})$  $(m_1 + m_2)\ddot{x}_1 + a_1\dot{x}_1 + a_2x_1 = u$ 

$$m_2(\ddot{x}_1 - \ddot{x}_2) = 0$$

• (ii) mode:  $m_2$ -body slides over  $m_1$ -body ( $\dot{x}_1 \neq \dot{x}_2$ )

$$m_1 \ddot{x}_1 + a_1 \dot{x}_1 + a_2 x_1 + b \operatorname{sgn}(\dot{z}) = u$$
  
$$m_2 \ddot{x}_2 - b \operatorname{sgn}(\dot{z}) = 0$$

- switching from mode (i) to mode (ii): if  $|\ddot{x}_1| > bm_2^{-1} \implies \ddot{x}_2 \neq \ddot{x}_1$
- switching from mode (ii) to mode (i): once  $\dot{x}_1 = \dot{x}_2 \implies$  synchronization



• Free system (2)-(4) without viscous damping (i.e.  $a_1=0$ ) is equivalent to

$$\dot{x}_{1} = v_{1}$$
(5)  

$$m_{1}\dot{v}_{1} = -a_{2}x_{1} - b\operatorname{sign}(v_{1} - v_{2})$$
(6)  

$$m_{2}\dot{v}_{2} = b\operatorname{sign}(v_{1} - v_{2})$$
(7)

Assuming Lyapunov function candidate, which is energy function for (5)-(7)

$$E = \frac{1}{2}a_2x_1^2 + \frac{1}{2}m_1v_1^2 + \frac{1}{2}m_2v_2^2$$

one obtains

 $\dot{E} = a_2 x_1 v_1 + m_1 v_1 \dot{v}_1 + m_2 v_2 \dot{v}_2 = \dots \dots = -b |v_1 - v_2| \le 0$ 

• System dissipates energy if  $m_2$  slides over  $m_1$ , i.e.  $v_1 \neq v_2 \implies \dot{E} < 0$ and it conserves energy if  $m_2$  rests on  $m_1$ , i.e.  $v_1 = v_2 \implies \dot{E} = 0$ 

Convergence



- For the intervals with *E*<0, that implies the system (2)-(4) is (locally) asymptotically stable on the intervals, consider the unforced dynamics of (3)  $m_1\ddot{x}_1 + a_1\dot{x}_1 + a_2x_1 = \pm b$
- One can show that

$$\ddot{x}_{1}(t) = \frac{\mathrm{d}^{2}}{\mathrm{d}t^{2}} \left[ \pm b \left( 1 - \frac{1}{\sqrt{1 - \zeta^{2}}} \exp(-\zeta \omega_{n} t) \sin\left(\omega_{n} \sqrt{1 - \zeta^{2}} t + \phi\right) \right) \right]$$
  
where  $\omega_{n} = \sqrt{\frac{a_{2}}{m_{1}}}, \quad \zeta = \frac{a_{1}}{2m_{1}\omega_{n}}, \quad \phi = \tan^{-1} \left( \frac{\sqrt{1 - \zeta^{2}}}{\zeta} \right)$ 

is continuously decaying on all intervals where E<0, that is in (ii)-mode

• Therefore, at some  $T_s < t < \infty$  $|\ddot{x}_1(t)| < b m_2^{-1} \implies \text{eq. (4) is no longer switching, i.e. (i)-mode perstains}$  $\implies \dot{x}_1(t) = \dot{x}_2(t) \quad \forall t > T_s$ 





[Ruderman, Zagvozdkin, Rachinskii, IEEE CDC, 2022]







# System with frictional interface and impact





| sampling rate                 | 2 kHz |
|-------------------------------|-------|
| laser sensor<br>repeatability | 8 µm  |

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# Some observations from impact experiments







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#### Problems in modeling vibroimpact



FORCES ONLY COMPRES -SIVE

(b)



(c) FORCES ARE COMPRESSIVE (x POSITIVE) AND TENSILE (x NEGATIVE)

[Hunt, Crossley, Journal of App. Mechanics, 1975]

Spurious jumps in the contact force if allowing for a linear damping term



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### Problems in modeling vibroimpact (cont.)

#### **Overall nonlinear contact force**

Closed hysteresis cycle of development of the contact force (here different impact velocities  $v_i$ )

[Hunt, Crossley, Journal of App. Mechanics, 1975]

$$f = -b\dot{x} - kx$$











 $f = -(\lambda z^n)\dot{z} - kz^n \qquad \lambda = 1.5\alpha k$ 



Parameter of the impacting surfaces n. For most simple assumption of the flat impacting surfaces n=1, while n=1.5 is consistent with Herzian theory of colliding spheres





#### Experimental results

# Hammer impact experiments





 $x_2$ 

Constant velocity of the platform drive experiments (i.e.  $dx_1/dt = const$ )





## Previous referred works

- 1. Dynamics of inertial pair coupled via frictional interface. M. Ruderman, A. Zagvozdkin, D. Rachinskii. IEEE 61th Conference on Decision and Control (CDC), 2022, pp. 1324-1328
- 2. On stiffness and damping of vibro-impact dynamics of backlash. M Ruderman. IEEE 30th International Symposium on Industrial Electronics (ISIE), 2021, p. 4
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- Analysis and compensation of kinetic friction in robotic and mechatronic control systems. M. Ruderman, CRC Press, 1<sup>st</sup> edition, 2023
- 5. Coefficient of restitution interpreted as damping in vibroimpact. K. Hunt, F. Crossley. Journal of Applied Mechanics, vol. 42 (2), 1975, pp. 440-445



# Thank you for attention

# Grazie per l'attenzione

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